

This article was downloaded by: [Tomsk State University of Control Systems and Radio]

On: 20 February 2013, At: 12:05

Publisher: Taylor & Francis

Informa Ltd Registered in England and Wales Registered Number: 1072954  
Registered office: Mortimer House, 37-41 Mortimer Street, London W1T 3JH, UK



## Molecular Crystals and Liquid Crystals

Publication details, including instructions for authors and subscription information:

<http://www.tandfonline.com/loi/gmcl16>

### Electrical Response Of $K_{0.30}MoO_3$ Around Incommensurate-Nearly Commensurate Change At 100k

Kitomi Tsutsumi <sup>a</sup>, Tsuyoshi Tamegai <sup>a</sup>, Seiichi Kagoshima <sup>a</sup> & Masatoshi Sato <sup>b</sup>

<sup>a</sup> Department of Pure and Applied Sciences,  
University of Tokyo, Komaba 3-8-1, Meguro, Tokyo,  
153, Japan

<sup>b</sup> Institute for Solid State Physics, University of  
Tokyo, Roppongi 7-22-1, Minato, Tokyo, 106, Japan  
Version of record first published: 20 Apr 2011.

To cite this article: Kitomi Tsutsumi, Tsuyoshi Tamegai, Seiichi Kagoshima & Masatoshi Sato (1985): Electrical Response Of  $K_{0.30}MoO_3$  Around Incommensurate-Nearly Commensurate Change At 100k, Molecular Crystals and Liquid Crystals, 121:1-4, 129-132

To link to this article: <http://dx.doi.org/10.1080/00268948508074846>

PLEASE SCROLL DOWN FOR ARTICLE

Full terms and conditions of use: <http://www.tandfonline.com/page/terms-and-conditions>

This article may be used for research, teaching, and private study purposes. Any substantial or systematic reproduction, redistribution, reselling, loan,

sub-licensing, systematic supply, or distribution in any form to anyone is expressly forbidden.

The publisher does not give any warranty express or implied or make any representation that the contents will be complete or accurate or up to date. The accuracy of any instructions, formulae, and drug doses should be independently verified with primary sources. The publisher shall not be liable for any loss, actions, claims, proceedings, demand, or costs or damages whatsoever or howsoever caused arising directly or indirectly in connection with or arising out of the use of this material.

## ELECTRICAL RESPONSE OF $K_{0.30}MoO_3$ AROUND INCOMMENSURATE-NEARLY COMMENSURATE CHANGE AT 100K

KITOMI TSUTSUMI, TSUYOSHI TAMEGAI, SEIICHI KAGOSHIMA  
Department of Pure and Applied Sciences, University of Tokyo,  
Komaba 3-8-1, Meguro, Tokyo 153, Japan

MASATOSHI SATO  
Institute for Solid State Physics, University of Tokyo,  
Roppongi 7-22-1, Minato, Tokyo 106, Japan

Abstract Electrical response and structural change were investigated in  $K_{0.30}MoO_3$  in the charge-density wave state around 100K, where the wave vector of the charge-density wave stops changing.

### INTRODUCTION

Non-linear transport properties are the subjects of recent investigations of one-dimensional conductors. Non-linear conductivity<sup>1</sup>, narrow band noise<sup>2</sup>, and broad band noise are attributed to the depinning of the charge-density wave (CDW) by electric fields. Electrical hysteresis<sup>3</sup> and the transient response observed in the trichalcogenides<sup>4</sup> and blue bronzes<sup>5</sup> are considered to be related to the metastability of CDW. In this article, we report the electrical response and the structural change of  $K_{0.30}MoO_3$ .

### EXPERIMENTAL RESULTS AND DISCUSSIONS

As far as the electrical response is concerned, we found two types of samples of  $K_{0.30}MoO_3$ . One is the "switching" sample which shows sudden onset of non-linear conduction with a resistance drop at the threshold field. The other is the "non-switching" sample which smoothly decreases the resistance above the threshold field as in  $NbSe_3$ . For bidirectional current pulses exceeding the threshold

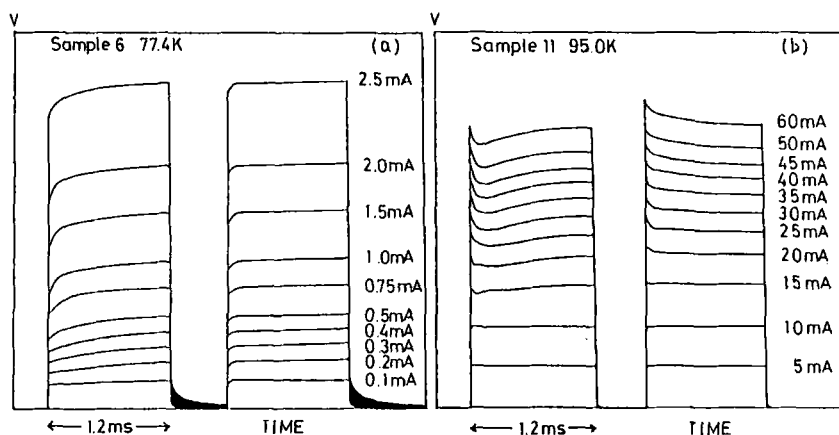


FIGURE 1 Responses to positive current double pulses after applying enough negative current pulse. (a) is "non-switching" sample ( $I_C=0.10\text{mA}$ ). (b) is "switching" sample ( $I_C=16\text{mA}$ ) current, non-switching samples show one type of transient response as shown in Fig. 1(a). Just after the application of the current pulse, the sample is in more conductive state. Then it relaxes to the steady state. Switching sample show another kind of relaxation as shown in Fig. 1(b). The cusp is considered to be related with the pinning of CDW. Both kinds of relaxations decrease their magnitude above about 100K. For both kinds of samples threshold field increases with increasing temperature.

To study the structural change due to the field, time resolved x-ray diffraction experiments were made under electric fields using a synchrotron radiation facility (SOR) in KEK at Tsukuba. X-ray wave length of  $1.198\text{\AA}$  was used. Diffraction intensity was measured by a solid-state detector. Practical resolution estimated from the widths of ordinary Bragg reflections are  $0.009\text{\AA}^{-1}$ ,  $0.003\text{\AA}^{-1}$ , and  $0.009\text{\AA}^{-1}$  along the  $b^*$ -,  $(2a^*-c^*)$ -, and  $(a^*+2c^*)$ -directions, respectively. Bidirectional pulsed electric currents were applied. Detected x-ray signals were accumulated by a multichannel analyzer in the

multichannel scaling (MCS) mode. The diffracted beam intensity is time dependent indicating a movement of the satellite position in the reciprocal space (i.e. the change of the wave vector of CDW-ordering). By measuring the profile of the satellite reflection, we found the temporal evolution of the satellite reflection occurs in the  $(2a^*-c^*)$ -direction. Figure 2 shows the the peak position and the width of the satellite. One finds two stable positons for the

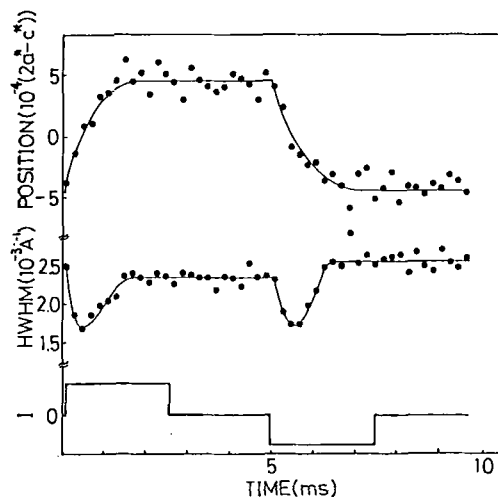


FIGURE 2 The position and the HWHM (half width at half maximum) of the satellite as a function of bidirectional current pulse.

positive and the negative polarities, respectively. These stable states are also seen by applying d.c. currents. The transient time from one state to the other is of the order of 1 ms, which seems to correspond to the electrical transient time. In the transient period the correlation length along the  $(2a^*-c^*)$ -direction deduced from the width of the satellite is longer than the detectable limit. In the steady state, however, it is estimated to be  $\sim 500\text{\AA}$ . The change in the CDW-ordering described above is most remarkably seen around 70K. Above 90K and below 40K no such change is observed.

The change of the wave vector in the direction of  $(2a^*-c^*)$

indicates that the phase difference between CDW in the neighboring infinite layers of  $\text{MoO}_6$  varies as a function of electric fields. Figure 3 depicts a possible configuration of CDW in real space. These two configurations may be caused by a field gradient in the sample. Field gradient results in the gradient of the force exerted to CDW by the field. Lack of the change in CDW-ordering above 90K may correspond to the decrease of the magnitude of the transient electrical relaxation above 100K.

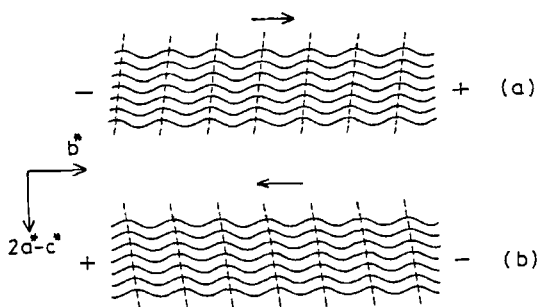


FIGURE 3 Possible configuration of CDW in real space. (a),(b) Each sinusoidal curve represents CDW in each infinite layer of  $\text{MoO}_6$ . Dotted lines show planes of the same phase. Arrows indicate the directions of the sliding of CDW.

#### ACKNOWLEDGEMENTS

The authors are very grateful for the assistance of the working group for low temperature diffraction in KEK. They also acknowledge Professors Y. Yamada and H. Fukuyama for stimulating discussions.

#### REFERENCES

1. P. Monceau, N. P. Ong, A. M. Portis, A. Meerschaut, and J. Rouxel, *Phys. Rev. Lett.*, **37**, 602 (1976).
2. R. M. Fleming and C. C. Grimes, *Phys. Rev. Lett.*, **42**, 1432 (1979).
3. A. W. Higgs and J. C. Gill, *Solid State Commun.*, **47** 737 (1983).
4. R. M. Fleming, *Solid State Commun.*, **43**, 167 (1982).
5. R. M. Fleming and L. F. Schneemeyer, *Phys. Rev. B* **28**, 6996 (1983)